

ter (50:30:14:6). In both cases the development was started immediately after the addition of solvent, and without tank saturation.

For the detection of the amino acids on the chromatogram a ninhydrin-collidine reagent was used. The reagent consisted of ninhydrin (0.15 g), acetic acid (10 ml) and 2,4,6-trimethyl-pyridine (1.25 ml), made up to a final volume of 50 ml with ethanol.² After spraying, the plates were placed on a hot plate (125°C) and the shape and colour of the spots marked with a pencil as they appeared. Not only was identification possible from the positions of the spots, but the process was also made easier by noting the speed of development, intensity and shade.

In addition to the 23 amino acids resolved by Jones and Heathcote, the system also provides a separation for norleucine, β -alanine, 4-aminobutyric acid, and 2-aminoglucose, some of which are present in complex humic hydrolysates. Norleucine (9.3, 5.9) is positioned close to, but definitely separated from, leucine (9.2, 5.4) and isoleucine (8.9, 4.8), where the parentheses give the R_F values for the first and second dimensions. β -Alanine (5.0, 1.1) and α -alanine (6.1, 1.3) separate well, but the latter may be partly overlapped by 4-aminobutyric acid (5.8, 1.2) which is positioned between them. Nevertheless all three can be recognized when present together in about equal amounts. 2-Aminoglucose (2.9, 3.3) is situated near serine (3.3, 1.8), but well separated from it.

According to the technique outlined above, we have analysed samples arising from different treatments of raw humus. The system works satisfactorily and also resolves several ninhydrin-positive spots which could not be assigned to any of the compounds mentioned above.

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Redetermined Crystal Structure of FeS₂ (Pyrite)

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The crystal structure of pyrite (FeS₂) was one of the very first to be determined by the method of X-ray diffraction. This early (1914) work by Bragg¹ and Ewald and Friedrich² depended on the use of single crystals, as are required by the Laue method. Suitable crystals were available in the form of the mineral pyrite. The only redetermination of this structure on the basis of an improved single crystal technique was reported in 1932 by Parker and Whitehouse,³ whereas other investigators have employed various powder techniques on natural as well as synthetic samples.

Chemical intuition and elementary semi-theoretical considerations led Pauling and Huggins⁴ (see also Pauling⁵) to suggest that the shortest (*i.e.* bonding) interatomic Fe—S and S—S distances in the pyrite and marcasite modifications of FeS₂, respectively, should be virtually equal. These expectations were found⁴ to be confirmed with respect to the bonding Fe—S distances (which at that time were reported to be identical), whereas the short S—S pair distance appeared to differ considerably (2.09—2.14 Å *versus* 2.25 Å, respectively) between the two modifications. According to their empirical scheme of tetrahedral radii, Pauling and Huggins⁴ furthermore deduced a predicted S—S distance of 2.08 Å, which together with the above evidence led them to propose that the parameter values for marcasite reported by Buerger⁶ were probably inaccurate. Provoked by this criticism,^{4,7} Buerger⁸ repeated the determination of the marcasite structure and confirmed his original⁶ parameter values with improved accuracy. However, the question concerning the distinction between the S—S pair distance in pyrite and marcasite remained unresolved, since the positional parameter of the pyrite structure might equally well be inaccurate. This view is in fact supported by the scattered values reported over the years for the S—S pair distance in pyrite (ranging from 2.09 to 2.17 Å, the latter value by Elliott⁹ being the most

recent), thus leaving the value uncertain, even up to the present. It therefore appeared to be of considerable interest to attempt a definitive redetermination of the structure by a single crystal technique and the use of modern computing and correction methods.

Single crystals of FeS_2 -*p* were prepared by chemical transport reactions, using iodine (~ 1 mg/ml capsule volume) as the transport agent. A good yield of suitable crystals was obtained by applying a temperature gradient of $\sim 1^\circ\text{C}/\text{mm}$ along a ~ 200 mm long (evacuated and sealed) silica capsule, with the hot end containing a stoichiometric mixture of the elements (99.99+ % Fe from Johnson, Matthey & Co., Ltd. and 99.9999 % S from Koch-Light Laboratories, Ltd.), at 620°C . These conditions produced a considerable number of crystals at the cold end of the capsule after 5 days.

The majority of the crystals were shaped as rectangular prisms, which showed a large variation in proportions. The crystal selected for the present study had the dimensions $0.07 \times 0.04 \times 0.03$ mm, the crystallographic axes being parallel to the edges of the crystal. X-Ray photographs were taken in an integrating Weissenberg camera of 57.3 mm diameter with $\text{MoK}\alpha$ -radiation using the multiple-film technique and the intensities were measured microphotometrically except for the weakest reflections which were estimated visually. The intensities were corrected for the combined Lorentz and polarization factors, and for absorption and secondary extinction according to the actual shape of the crystal. Atomic scattering factors were taken from Hanson *et al.*¹⁰ and the corrections and refinements were carried out according to programmes by Dahl *et al.*¹¹

The dimension $a = 5.4179 \pm 0.0011 \text{ \AA}$ * of the cubic unit cell of FeS_2 -*p* (deduced from Guinier photographic data of crushed crystals according to the method of least squares) is in perfect agreement with the value recently reported by Bither *et al.*¹² for a synthetic sample. FeS_2 -*p* assumes the stoichiometric 1:2 composition without having any appreciable range of homogeneity and the unit cell contains four formula units. The systematic extinctions of reflections of the type $0kl$ for $k=2n+1$ ($h0l$ for $l=2n+1$ and $hk0$ for $h=2n+1$)

* Throughout this paper the indicated error limits correspond to the standard deviations.

in the diffraction data, imply that the space group is $Pa\bar{3}$.

The pyrite structure has been confirmed in several independent studies (*cf.* Refs. 1–3,9) and its essential features are clearly correct. Accepting that the space group is precisely established from the systematic extinctions, there are one positional parameter and two temperature factors to be determined for this structure.** The least squares refinements converged very rapidly to the following parameter values: Fe in 4(*a*) with $B = 0.14 \pm 0.04 \text{ \AA}^2$; S in 8(*c*) with $x = 0.3840 \pm 0.0005$, $B = 0.12 \pm 0.05 \text{ \AA}^2$, with $R = \frac{\sum |F_o| - |F_c|}{\sum |F_o|} = 0.075$. (The unobserved reflections are not included in the least squares refinements and the *R*-value.)

The present value for the positional parameter is identical with that derived by Elliott⁹ from powder data. Despite the fact that Elliott was unable to use an objective criterion in assessing the probable error in the determination, it is worth noting that his error limits correspond to twice the standard deviation obtained in the present study.

The present values for *a* and *x* lead to the following bonding interatomic distances and relevant angles in pyrite:

Fe–S(6)	$2.262 \pm 0.003 \text{ \AA}$
S–S(1)	$2.177 \pm 0.004 \text{ \AA}$
S–Fe–S(6+6)	$85.57 \pm 0.10^\circ$
	$94.43 \pm 0.10^\circ$
Fe–S–Fe(3)	$115.71 \pm 0.12^\circ$
S–S–Fe(3)	$102.13 \pm 0.13^\circ$

The length of these distances may be compared with the corresponding values obtained¹⁴ for marcasite:

Fe–S (2+2+2)	2.230 ± 0.006 ,
	2.235 ± 0.002 ,
	$2.275 \pm 0.006 \text{ \AA}$;
	average 2.247 \AA
S–S(1)	$2.223 \pm 0.003 \text{ \AA}$

It is seen that both the bonding Fe–S and S–S distances differ significantly between the two modifications, although the difference between the S–S distances in the two structures is rather less than was

** This neglects the possibility of infinitesimal deviations from the symmetry of space group $Pa\bar{3}$. A reduction of the symmetry to that of space group $P2_13$ would, for example, provide virtually the same atomic arrangement, with a slight modification of the pyrite structure towards that of ullmanite (NiSbS).¹³

considered to be the case at the time of the controversy between Pauling and Huggins^{4,7} on the one hand and Buerger^{6,8} on the other (*vide supra*). These differences between the two modifications, which are related to the structural dissimilarities, will be subject to discussion in a forthcoming paper.

As a consequence of the cubic symmetry of the pyrite type structure and the fact that there is only one positional parameter, knowledge of the lattice constant a and estimates of the length of the bonding $X-X$ ($a\sqrt{3}(1-2x)$) or $T-X$ ($a(3x^2-2x+\frac{1}{2})^{\frac{1}{2}}$) distance enables prediction of the value of x (*cf.*, *e.g.*, Pauling and Huggins⁴). More conveniently, however, the value of x may be estimated from the ratio between these distances ($\sqrt{3}(1-2x)/(3x^2-2x+\frac{1}{2})^{\frac{1}{2}}$), since it may be easier to judge this quantity with reliability. (In the relevant range of x , *i.e.* $0.35 < x < 0.42$, the ratio can to a good approximation be simplified to the linear function $4.4454 - 9.0722x$.)

Substitution of the observed distance ratio for the marcasite modification of FeS_2 in the above expression, gives $x=0.381$ for pyrite, which accords reasonably well with the experimental value 0.3840. Similarly, the recent values for the distance ratios in the marcasite modifications of FeTe_2 ,¹⁴ CoTe_2 ,¹⁴ and $\beta\text{-NiAs}_2$,¹⁵ may be utilized for the prediction of the x -values 0.364, 0.367, and 0.377, respectively, for the corresponding modifications with the pyrite type structure which have recently been prepared^{12,16} by high pressure syntheses. The hitherto unobserved pyrite variant of NiSb_2 may be predicted to have $a=6.204$ Å and $x=0.367$ on the basis of the structural data¹⁶ for the marcasite modification. It must be emphasized, however, that these predictions are subject to some uncertainty (in particular in the latter case) since, *e.g.*, the size of transition metal atoms appears to correlate with their localized configurations of unpaired electrons (*cf.* Pearson¹⁷).

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New Metal-rich Arsenides of Niobium and Tantalum

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In several earlier investigations of the Nb-As and Ta-As systems,¹⁻⁷ the method of synthesis has been the direct reaction between the component elements at temperatures up to 1000°C. Under these conditions the only intermediate phases formed in each system are a monoarsenide and a diarsenide. Ganglberger,^{8,9} who employed higher temperatures for his syntheses, was able to prepare the metal-rich compounds Nb_3As and Ta_3As , both of which were reported to crystallize with the Ti_3P -type structure.

In the present paper, some results are reported from an examination of the

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