# On the Phase Relationships in the Titanium-Antimony System. The Crystal Structures of Ti<sub>2</sub>Sb

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By using different melting methods four new titanium antimonides have been identified:

The cubic  $Ti_3Sb$  phase, prepared in an arc furnace, is found to have the lattice constant a=5.2186 Å. The pycnometric density is 6.137 g cm<sup>-3</sup> at 25°C. The crystal structure of the cubic  $Ti_3Sb$  phase is of the  $\beta$ -W type.

The tetragonal Ti<sub>3</sub>Sb phase, prepared in alumina crucibles in a carbon resistance furnace, is found to have a body-centered unit cell with dimensions: a=10.465 Å, c=5.2639 Å. The pycnometric density is 6.070 g cm<sup>-3</sup> at 25°C. The space group is I4/mcm  $(D_{4b}^{18})$  and the unit cell contains 24 Ti and 8 Sb in the positions: (0,0,0; 1/2,1/2,1/2) + 4 Ti in  $(a) \pm 0,0,1/4; 4$  Ti in (b) 0,1/2,1/4;

1/2,0,1/4; 16 Ti in  $(k) \pm (x,y,0;\overline{y},x,0;x,\overline{y},1/2;y,x,1/2)$  with  $x_1=0.0766,$   $y_1=0.2228$  and 8 Sb in  $(h) \pm (x,1/2+x,0;1/2+x,\overline{x},0)$  with  $x_1 = 0.1635$ .

The structure of the tetragonal Ti<sub>3</sub>Sb phase is closely related to the

β-W type structure of the cubic Ti<sub>3</sub>Sb phase.

Cubic crystals are transformed into tetragonal crystals on melting in alumina crucibles, while tetragonal crystals remain tetragonal after melting in the arc furnace. Possible reasons for the apparently irreversible transformation of the two phases are discussed.

The orthorhombic Ti<sub>1.7</sub>Sb phase, prepared in the carbon resistance furnace, is found to have lattice constants a = 10.172 Å, b = 8.348 Åand c = 7.135 Å. The probable space groups are Pmma  $(D_{2k}^5)$ ,  $Pmc2_1$  $(C_{2v}^2)$  and  $Pma2\ (C_{2v}^4)$ .

The orthorhombic Ti<sub>1.2</sub>Sb phase, prepared in the carbon resistance furnace, is found to have lattice constants a = 14.55 Å, b = 16.34 Å,and c = 5.31 Å. The probable space groups are Pbam  $(D_{2h}^9)$  and  $Pba2 \ (C_{2v}^8).$ 

The redetermined lattice constants of the hexagonal TiSb phase, prepared in the carbon resistance furnace, are  $\ddot{a} = 4.115$  Å and c = 6.264 Å.

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As result of X-ray studies of titanium antimony alloys Nowotny and cowor-Akers  $^{1-3}$  reported the existence of four intermediate phases:  $\mathrm{Ti_4Sb}$ ,  $\mathrm{Ti_{2.5}Sb}$ ,  $\mathrm{TiSb}$  and  $\mathrm{TiSb_2}$ . The  $\mathrm{Ti_4Sb}$  phase  $^{1,2}$  was prepared by heating weighed amounts of pure antimony and Degussa titanium in porcelain (Pythagoras) tubes in a carbon resistance (Nernst-Tammann) furnace at  $1600^{\circ}\mathrm{C}$ . The alloys were afterwards annealed at  $600^{\circ}\mathrm{C}$  and quenched, or in some cases slowly cooled, to room temperature. The crystallographic properties of this phase were determined from X-ray powder photographs. By comparing visually estimated intensities with those obtained by calculations, Nowotny et al. concluded that the  $\mathrm{Ti_4Sb}$  phase had a structure of the  $\mathrm{Ni_3Sn}$ -( $\mathrm{D0_{19}}$ )-like type. The deviating composition was accounted for by assuming 0.2 titanium atoms, corresponding to the formula  $\mathrm{Ti_3(Ti_{0.2}Sb_{0.8})}$ , to be distributed at random over the antimony positions.

The  $Ti_{2.5}Sb$  phase 3 was synthesized by melting appropriate amounts of the elements in crucibles of carbon using either a Tammann high pressure furnace or a high frequency induction furnace. The crystal structure of  $Ti_{2.5}Sb$  was determined by Auer-Welsbach *et al.*³ by means of single crystal and powder photographs to be of the  $Cu_2Sb-(C38)$ -like type but with doubled *c*-axis. The parameter values  $z_{Ti}=0.133$  and  $z_{Sb}=0.353$  were determined from the intensities of the 00l reflections. Substitution of 0.15 antimony atoms for titanium atoms, corresponding to the formula  $Ti_2(Ti_{0.15}Sb_{0.85})$  was thought to

explain the excess of titanium and the observed density.

The TiSb phase with NiAs-(B8)-type structure and the TiSb<sub>2</sub> phase with CuAl<sub>2</sub>-(C16)-type structure were prepared and examined <sup>1,2</sup> as described for the Ti<sub>4</sub>Sb phase. These two phases were also obtained by melting titanium and antimony in the right proportions in open crucibles under halide flux.

Using different methods of preparation, the existence of four different, intermediate phases has been established in this institute <sup>4,5</sup>. The present paper gives an account of the crystal structure determinations for two different phases with composition Ti<sub>3</sub>Sb. In addition, some preliminary phase-analytical observations of two more titanium antimonides are reported.

After being informed about the present work, Matthias et al.<sup>6</sup> have duplicated the preparation of the cubic  $Ti_3Sb$  phase and shown that it is a supercon-

ductor with a transition temperature of 5.8°K.

### **EXPERIMENTAL**

The titanium metal used in this study was "iodide" titanium in rods supplied through A. D. MacKay, Inc., New York. A spectrographic analysis showed the presence of about 30 ppm of Fe and Mn and less than 10 ppm Cu. No analysis of non-metallic impurities was carried out.

The antimony was "High purity antimony" from Bradley Mining Co., San Francisco. According to the supplied analysis it contained the following impurities (in ppm): As

(10), Cu (3), Fe (3) and Pb (2).

Several different techniques were tried for preparing the titanium-antimony alloys. In some preliminary experiments alloys were prepared by heating titanium and antimony in evacuated, sealed silica tubes. However, contamination of the alloys due to reaction with silica then always took place. According to Auer-Welsbach *et al.*<sup>3</sup> the compound  $\mathrm{Ti}_5\mathrm{Si}_3$  is formed by this reaction.

One series of alloys with composition ranging from Ti<sub>6</sub>Sb to TiSb was prepared in an electric arc furnace (Degussa Volio) by melting weighed amounts of titanium and antimony, previously mixed and compacted. A protective atmosphere of argon gas was present in the furnace during the melting process. Using a pressure of about 4 atm no preferential loss of antimony from the samples was seen, while melting in vacuo gave visible loss of antimony.

In the same composition range a second series of alloys was made by melting weighed quantities of the elements in crucibles of pure alumina, placed inside silica tubes which were evacuated and sealed. The samples were melted in a carbon resistance (Nernst-Tammann) furnace at temperatures ranging from 1400 to 1600°C. Temperatures were measured with an optical pyrometer. When the furnace reached the desired temperature, the sample was immediately removed and quenched in ice water. If the temperature had exceeded 1550°C, the silica tubes often broke during quenching.

The alloys could not be heated in contact with alumina at these temperatures for a longer time, as reactions then were seen to take place. Furthermore it was noticed that pure titanium is able to reduce alumina. This reaction was unsuspected in view of the thermochemical discussion by Kubaschewski and Dench 7 related to the production of Ti-Al alloys from TiO2 and excess Al at 1000°C. Nevertheless, a reduction of alumina

might apparently well take place at higher temperatures

The melted alloys were bright shiny and metallic looking, while the sintered alloys were in form of a gray to black powder. The alloys were crushed and finely ground and X-ray powder photographs were taken in cameras with 114.6 mm effective diameter and X-ray powder photographs were taken in cameras with 114.0 influence to diameter and asymmetric film mounting. In most cases filtered copper radiation was used, while chromium radiation was used to obtain greater dispersion in the photographs of the tetragonal Ti<sub>3</sub>Sb samples. The lattice constants are expressed in Å units, taking  $\lambda_{\text{Cu}K\alpha_1} = 1.54050$  Å and  $\lambda_{\text{Cr}K\alpha_1} = 2.28962$  Å. Results from Nowotny and coworkers <sup>1-3</sup>, expressed in kX, have been transformed to Å by multiplication by the factor 1.00202. The lattice constant of the cubic phase was found by extrapolation according to the method of Nelson and Bilays. to the method of Nelson and Riley 8.

Oscillation and Weissenberg photographs were taken with CuK-radiation using a Weissenberg camera of 57.3 mm diameter.

The density of the samples was determined pycnometrically at  $25^{\circ}$ C with kerosene as displacement liquid. To remove gases adsorbed by the sample, the pycnometer was filled with kerosene under vacuum. The samples weighed from 2 to 3 g.

### RESULTS

In the titanium-antimony system the existence of four new, intermediate phases was established, i.e. the cubic and tetragonal Ti<sub>3</sub>Sb phases, the Ti<sub>1.7</sub>Sb phase and the Ti<sub>1,2</sub>Sb phase. Furthermore, the existence of the TiSb phase reported by Nowotny and Pesl was confirmed. The TiSb<sub>2</sub> phase found by Nowotny and Pesl was not studied here. By the experimental technique used in this study no samples containing the Ti<sub>4</sub>Sb and Ti<sub>2.5</sub>Sb phases reported by Nowotny and coworkers 1-3 were obtained.

## A. The cubic Ti<sub>3</sub>Sb phase

Powder photographs of samples prepared in the arc furnace showed the presence of a phase which could be indexed on the basis of a cubic unit cell, cf. Table 1. Inferred from the powder photographs, the amount of the other phases seems to disappear close to the stoichiometric ratio 3:1. A series of density measurements, cf. Table 2, was then carried out for an accurate determination of the composition of this phase. The specific volume versus weight % Sb plot shown in Fig. 1 had a break in the linearity at 45.9 weight % Sb (25.0 atomic % Sb), ascertaining the composition to be Ti<sub>3</sub>Sb.

Table 1.	Powder	photograph	data	of	cubic	Ti <sub>3</sub> Sb	taken	with	CuK-r	adiation,	$a_2$	$\mathbf{and}$	β
						mitted							

$\sin^2\Theta \times 10^4$	hkl	$I_{ m obs}$	$I_{ m calc}$	$\sin^2\!\Theta \times 10^4$	hkl	$I_{ m obs}$	$I_{\mathrm{calc}}$
_	100	0	0	5670 {	510	2.4	0.8
437	110	12.3	14.3	3070	431	] 2.4	1.5
_	111	0	0	_	511	o i	0
873	200	13.9	15.5	_	333	} • {	0
1092	210	14.6	13.8	6330	520	2.5	0.9
1312	211	30.2	34.1		<b>432</b>	1) (1	1.8
1746	220	2.7	2.2	6538	521	6.9	6.6
	300	0 {	0	6978	<b>440</b>	3.6	3.9
_	221		0	_	$\boldsymbol{522}$	) o {	0
2181	310	3.8	3.1	_	441		0
	311	0	0	7412	530	1.8	0.8
-	<b>222</b>	0	0.02	1412	433	II II	0.8
2837	320	3.4	2.6		531	0	0
3056	321	17.7	15.8	7849.6	600	5.5	0.9
3491	400	4.2	3.8	l l	<b>442</b>	IJ (i	3.8
_	410	0 {	0	8065	610	1.3	1.1
-	$\bf 322$	1) 1	0	8284.5	611	12.8	4.0
3929	411	1 2.3 $1$	1.2	i ti	$\bf 532$	I) (I	8.2
1	330	<b>)</b>	0.6	8722	620	1.7	1.1
_	331	0	0	-	621	1) (1	0
4361	<b>420</b>	4.4	4.5	-	540	} 0 {	0
4582	421	2.8	2.4		<b>443</b>	, ,	0
4800	332	4.7	4.0	9154.0	541	3.4	2.7
5236	<b>422</b>	0.9	0.8		533	0	0
-	500	) o {	0	-	622	0	0.00
-	430	1 , 1	0	9804.5	$\begin{array}{c} 630 \\ 542 \end{array}$	5.6	$\frac{3.0}{6.0}$

The lattice constant of the sample with the composition Ti<sub>3</sub>Sb is found in Table 2 together with the lattice constants of the Ti<sub>3</sub>Sb-phase in the samples Ti<sub>4</sub>Sb, Ti<sub>3.5</sub>Sb, Ti<sub>2.5</sub>Sb and Ti<sub>2</sub>Sb. As the lattice constant is almost invariant for specimens with different initial proportions of the components, no composition range of the cubic Ti<sub>5</sub>Sb-phase is noticeable.

composition range of the cubic  ${\rm Ti_3Sb\text{-}phase}$  is noticeable. The pycnometric density of  ${\rm Ti_3Sb\text{-}phase}$  is noticeable. The pycnometric density of  ${\rm Ti_3Sb}$  at 25°C, cf. Table 2, is found to be in agreement with the density calculated from the X-ray data, 6.201 g cm<sup>-3</sup>, assuming two ( $Z_{\rm c}=1.98$ )  ${\rm Ti_3Sb\text{-}groups}$  per unit cell.

Table 2. Lattice constants and pycnometric densities of samples containing the cubic Ti<sub>3</sub>Sb phase.

Sample	a (Å)	Density (g cm <sup>-3</sup> )
Ti <sub>4</sub> Sb	5.2187	5.801
Ti <sub>8.5</sub> Sb	5.2188	5.956
Ti <sub>3</sub> Sb	5.2186	6.137
Ti <sub>2.5</sub> Sb	5.2184	6.251
Ti <sub>2</sub> Sb	5.2184	6.381

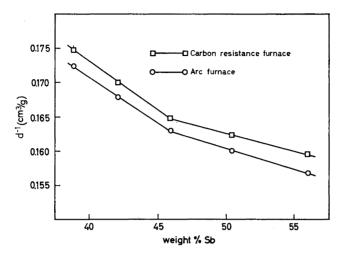


Fig. 1. Reciprocal densities as function of composition for samples synthesized in carbon resistance and arc furnace.

Many related Me<sub>3</sub>X compounds of the transition metals have a structure of the  $\beta$ -W (W<sub>3</sub>O) type. The  $\beta$ -W structure, cf. Fig. 2, is described in space group Pm3n ( $O_h^3$ ) with

6 Me in (d) 1/4,1/2,0; 0,1/4,1/2; 1/2,0,1/4; 3/4,1/2,0; 0,3/4,1/2; 1/2,0,3/4. 2 X in (a) 0,0,0; 1/2,1/2,1/2.

As the missing reflections in the powder photographs were in agreement with the extinctions for this space group, a comparison of observed and calculated

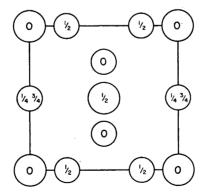


Fig. 2. Projection of one unit cell of the  $\beta$ -W type structure. Large circles represent antimony atoms and small circles titanium atoms. Figures give the z-parameters.

Table 3. Powder photograph	data of tetragonal Ti <sub>3</sub> Sb taken with CrK-radiation, a <sub>2</sub> an	d
• • • •	$\beta$ lines omitted.	

$\sin^2\Theta \times 10^4$	hkl	$I_{ m obs}$	$I_{ m calc}$	$\sin^2\Theta \times 10^4$	hkl	$I_{ m obs}$	$I_{ m calc}$
240	110	8.5	9.4	_	640	0	2.4
481	200	6.5	8.6	6298	413	46	39.2
	220	0	3.2	_	602	0	0.5
1074	211	40.5	38.1	6690	622	12	10.3
1200	310	14	16.4	6820	721	19.5	18.5
1898	002	35	33.7	6948	730	11.5	8.
_	400	0	0.0	_	433	0	0.
2034	<b>321</b>	31.5	26.1	7571	004	18	17.0
	112	0	1.3	_	800	0	0.3
2159	330	33.5	31.1	7731	523	35	28.
2378	202	50	48.7	_	114	0	0.4
2398	420	46.5	43.9		651	0	2.3
2514	411	113	124.8	7880	712	} 19 {	15.9
2853	$\boldsymbol{222}$	14.5	12.4	- 1	$\bf 552$	1) 1	5.'
3094	312	10.5	8.5		204	0	1.0
	510	0	0.0	8119	$\bf 642$	} 54.5 {	46.
_	431	0	0.0		820		12.
3814	<b>402</b>	18.5 {	14.4	8252	811	} 40.5 {	28.
-	<b>440</b>	D U	2.9		741	U U	8.
3949	521	43.5	38.5	8328	543	4.5	5.
4057	$\bf 332$	30.5	23.6	_	<b>224</b>	0	0.
	<b>53</b> 0	) 00.0 t	7.3	8613	660	32	30.
4299	<b>422</b>	} 17 {	12.9	_	613	0	0.
-	600	) - U	6.7	8765	314	3.5	6.
	620		0.3	8835	732	31.5	25.
4865	213	11.5 {	9.3		750	0	1.
-	611	)	0.1	9212	831	13.5	14.
7000	512	0	0.0		404	0	0.
5392	541	6.5	4.3	9550	802	33.5	31.
5001	442	0	0.9	9574	840	6	4.
5821	323	4 3	5.7	9643	633	11.5 35.5	12.
5868	$\begin{array}{c} 631 \\ 532 \end{array}$	5	4.5 10.0	9722	$\frac{334}{910}$	$\frac{35.5}{26.5}$	52.
5973	$\begin{array}{c} 532 \\ 710 \end{array}$	} <sub>26</sub> {		9813	910	20.0	38.
9819	550	Z0	$\begin{array}{c} 7.0 \\ 2.5 \end{array}$				
	000	<u> </u>	4.0				

intensities for the reflections was carried out. The intensities of the reflections were calculated using the formula

$$I_{
m calc} = 10^{-5} imes |F|^2 imes Lp imes 
u$$

For the calculation of the structure amplitudes, Thomas-Fermi model values of f, taken from Internationale Tabellen, were used throughout this paper. No corrections for absorption and temperature factors were applied. The observed intensities were obtained from photometer recordings and referred to a relative, absolute scale by making the sums of  $I_{\rm obs}$  and  $I_{\rm calc}$  equal. As can be seen from Table 1 the agreement between the two sets of values is very good, and leaves no doubt about the correctness of the proposed structure.

## B. The tetragonal Ti<sub>3</sub>Sb phase

Powder photographs of samples prepared in the Nernst-Tammann furnace show the presence of a phase with tetragonal structure, cf. Table 3. Powder photographs of various samples of Ti<sub>6</sub>Sb, Ti<sub>5</sub>Sb and Ti<sub>4</sub>Sb contain some of the strongest reflections from α-Ti, while powder photographs of Ti<sub>2</sub>Sb contain reflections from a phase richer in antimony with probable composition Ti<sub>1.7</sub>Sb. The additional reflections are weak on photographs of samples with composition Ti<sub>3.5</sub>Sb and Ti<sub>2.5</sub>Sb and are absent on photographs of samples with composition Ti<sub>3</sub>Sb. A more accurate determination of the composition of the phase was obtained from a plot of the specific volume versus weight % Sb, cf. Fig. 1, and by means of metallographic methods. Thus the composition of the phase is ascertained to be Ti<sub>3</sub>Sb.

The lattice constants of the tetragonal Ti<sub>3</sub>Sb phase in the samples Ti<sub>4</sub>Sb, Ti<sub>3.5</sub>Sb, Ti<sub>3</sub>Sb, Ti<sub>2.5</sub>Sb and Ti<sub>2</sub>Sb are listed in Table 4. The slight expansion of the unit cell, which takes place in the range Ti<sub>3</sub>Sb to Ti<sub>2.5</sub>Sb indicates that the tetragonal Ti<sub>3</sub>Sb phase has a narrow range of homogeneity on the antimony-rich side of Ti<sub>3</sub>Sb.

The observed density, cf. Table 4, indicates that the unit cell contains 8 ( $Z_c = 7.94$ ) Ti<sub>3</sub>Sb-groups. The calculated density for Ti<sub>3</sub>Sb from the X-ray measurements is 6.116 g cm<sup>-3</sup>.

Single crystals of  $Ti_3Sb$  were found in samples quenched from 1450°C. The single crystals were needle-shaped and oscillation and Weissenberg photographs showed tetragonal symmetry, Laue symmetry 4/mmm, with the c-axis along the needle axis. Missing reflections on the photographs were:

$$hkl$$
 when  $h+k+l=2n+1$   
 $hk0$  when  $(h+k=2n+1)$   
 $0kl$  when  $k=2n+1$  and  $l=2n+1$ 

Characteristic space groups are accordingly I4/mcm  $(D_{4k}^{18})$ , I4cm  $(C_{4v}^{10})$  and  $I\bar{4}c2$   $(D_{2d}^{10})$ . All possible space groups have symmetry p4m in the 001-projection with  $a'=b'=\frac{1}{2}\sqrt{2} a$ .

Determination of the structure. The intensities of the hk0 and hk1-reflections were estimated visually from "multiple film" Weissenberg photographs taken with CuK-radiation and were corrected for Lorentz and polarization factors. No corrections for absorption, extinction and temperature factors were applied.

Table 4. Lattice constants and pycnometric densities of samples containing the tetragonal Ti<sub>3</sub>Sb phase.

Sample	a (Å)	c (Å)	Density (g cm <sup>-3</sup> )
${f Ti_4Sb} \ {f Ti_3.5Sb} \ {f Ti_3Sb} \ {f Ti_{2.5}Sb} \ {f Ti_2Sb} \ {f Ti_2Sb}$	10.468	5.2640	5.721
	10.464	5.2637	5.876
	10.465	5.2639	6.070
	10.476	5.2688	6.161
	10.475	5.2692	6.270

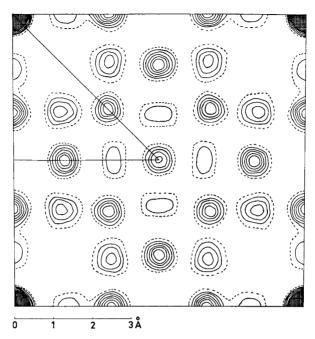


Fig. 3. Patterson projection on 001 of the tetragonal Ti<sub>3</sub>Sb. Only the primitive unit cell is drawn. The contours are at equal but arbitrary intervals. Zero contours are broken and negative contours are omitted.

On the basis of the corrected  $F(hk0)^2$ -values a Patterson map was computed, cf. Fig. 3. As the cubic and tetragonal Ti<sub>3</sub>Sb phase have related lattice dimensions  $a_{\rm tetr} \approx 2a_{\rm cub}$ ,  $c_{\rm tetr} \approx c_{\rm cub}$ , relationships with the  $\beta$ -W structure was looked for in the Patterson map. It was indeed possible to rearrange the atoms in the  $\beta$ -W structure in agreement with the Patterson map. Based on these approximate atomic parameters structure factors were calculated, and after adjusting  $F_{\rm o}$  and  $F_{\rm c}$  to the same scale  $R = \Sigma \big| |F_{\rm o}| - |F_{\rm c}| \big| / \Sigma |F_{\rm o}| = 0.29$  was found, indicating the main features of the proposed structure to be correct.

The atomic parameters were refined further by means of four Fourier and one difference synthesis. R decreased to 0.19 after the last Fourier synthesis and to 0.132 after the difference synthesis. The observed and calculated F(hk0)-values are listed in Table 5 and the final Fourier map is shown in Fig. 4. The corresponding parameters, referred to the body-centered axes, are:

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4 Ti I in 0,0 etc.

4 Ti II in 0,\frac{1}{2} etc.

16 Ti III in x_1 = 0.0766, y_1 = 0.2228 etc.

8 Sb in x_2 = 0.1635, y_2 = \frac{1}{2} + x_2 = 0.6635 etc.
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In terms of the highest symmetric space group, I4/mcm ( $D_{4h}^{18}$ ), all atoms are in special positions with fixed z-coordinates. For symmetry reasons only

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Table 5. Observed and calculated structure factors for tetragonal Ti<sub>3</sub>Sb (CuK-radiation).

hkl	$F_{\mathrm{o}}$	$\frac{1}{2}F_{ m c}$	hkl	$F_{ m o}$	$1/_2 F_{ m c}$
200	29.3	- 38.2	880	45.3	+ 37.8
400	0	- 1.4	1080	59.2	+ 51.3
600	110.7	+ 114.5			
800	16.0	- 23.9	990	105.7	-114.6
1000	25.6	- 22.9			
1200	152.5	+ 142.5	hkl	$F_{o}$	$1_{2} m{F}_{ m c} $
110	22.2	- 27.0		- 6	1 /2 - 6
310	70.3	+ 60.7	211	- 4 0	00.0
<b>51</b> 0	0	- 1.0	211	74.6	88.8
710	75.5	- 88.7	411	177.0	186.6
910	86.9	+ 91.9	611	0	6.4
1110	35.8	- 43.2	811	<b>95.2</b>	103.8
1310	0	- 5.7	1011	14.0	15.4
	40.0		1211	14.8	10.4
220	42.9	+ 33.2	321	68.0	71.4
420	139.1	+ 149.2	521		132.4
620	26.3	+ 18.1	721	118.4	77.4
820	101.9	+ 99.8	721	86.2	
1020	80.9	+ 70.8	921	0	3.8
1220	39.9	- 46.8	1121	104.4	102.8
			1321	59.0	70.0
330	141.5	-168.4	401	0	4.0
<b>53</b> 0	92.3	+ 82.8	431	0	
<b>73</b> 0	101.9	+ 92.9	631	48.4	50.2
930	<b>93.2</b>	- 96.5	831	54.6	64.0
1130	24.2	+ 16.6	1031	0	0.0
1330	51.9	+ 51.3	1231	45.4	44.0
440	85.5	+ 71.7	541	48.6	48.6
640	40.9	- 52.5	741	64.0	57.0
840	46.8	+ 37.5	941	41.0	37.2
1040	91.4	+ 76.2	1141	87.0	84.8
1240	0	+ 8.1			l
			651	37.8	31.4
550	69.5	+ 75.3	851	93.8	103.0
750	27.0	- 31.6	1051	134.6	114.6
950	0	+ 4.0	1251	9.0	11.0
1150	34.3	+ 25.4			
			761	11.8	6.4
660	209.9	+201.6	961	40.2	29.0
860	13.7	+ 8.3	1161	20.0	14.6
1060	28.8	- 36.4			
1260	105.2	+ 117.6	871	83.8	75.4
		1	1071	81.8	75.8
770	39.8	+ 30.5		·	1
970	46.4	+ 38.1	981	50.4	44.4
1170	0	- 3.9	1	į.	1

the 16 Ti III atoms in (k) and the 8 Sb atoms in (h) contribute to the hk1-reflections. The agreement between  $|F(hk1)|_{\circ}$  and  $|F(hk1)|_{\circ}$  calculated on this assumption, cf. Table 5, is so good, R=0.116, that the space group and atomic parameters  $x_1$ ,  $y_1$  and  $x_2$  seem correctly chosen.

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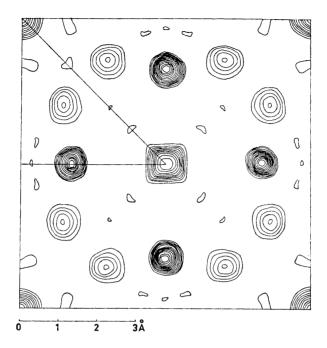


Fig. 4. Fourier projection of the primitive unit cell of tetragonal  $Ti_3Sb$  on 001. Contours are drawn at every 10 e Å<sup>-2</sup>. Zero and negative contours omitted.

On the basis of the relationships found between cubic and tetragonal  ${\rm Ti_3Sb}$  is seemed reasonable to assume that the z-coordinates in tetragonal  ${\rm Ti_3Sb}$  were the same as for the corresponding atoms in cubic  ${\rm Ti_3Sb}$ . It turned out, however, that z-values of 0 and  $\frac{1}{2}$  for Ti I were unreasonable considering the interatomic distances Ti I—Ti III. Therefore, z-parameter values of 1/4 and 3/4 were chosen for both Ti I and Ti II. In order to verify these z-parameters, intensities was calculated in the same way as for cubic  ${\rm Ti_3Sb}$  and compared with those obtained from photometer recordings of powder photographs. As can be seen from Table 3 the agreement between the two sets of values is good and leaves no doubt about the correctness of the proposed structure.

The deduced atomic arrangement, cf. Fig. 5, is thus: (0,0,0; 1/2,1/2,1/2) +

```
ted atomic arrangement, t/. Fig. 5, is thus: (0,0,0, 1/2,1/2,1/2) \neq 0

4 Ti I in (a) 0,0,1/4; 0,0,3/4.

4 Ti II in (b) 0,1/2,1/4; 1/2,0,1/4.

16 Ti III in (k) x,y,0; \overline{x},\overline{y},0; \overline{y},x,0; y,\overline{x},0; x,\overline{y},\frac{1}{2}; \overline{x},y,\frac{1}{2}; \overline{y},x,\frac{1}{2}; \overline{y},\overline{x},\frac{1}{2}, with <math>x_1 = 0.0766, y_1 = 0.2228.

8 Sb in (h) x,\frac{1}{2} + x,0; \overline{x},\frac{1}{2} - x,0; \frac{1}{2} + x,\overline{x},0, \frac{1}{2} - x,x,0 with <math>x_2 = 0.1635.
```

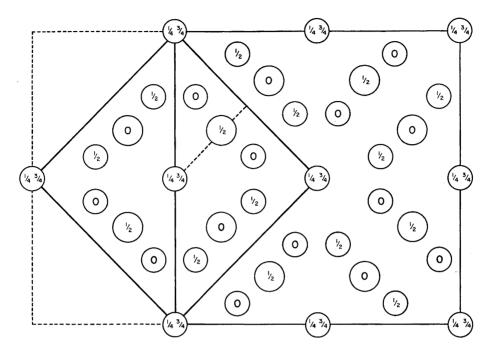


Fig. 5. Projection of the unit cell of tetragonal Ti<sub>3</sub>Sb on 001. Large circles represent antimony atoms and small circles titanium atoms. Figures give the z-parameters. Both the primitive and the body-centered unit cell are indicated.

### C. Some phase-analytical data

In addition to the two  $\rm Ti_3Sb$  phases just reported, two other previously unknown phases, the  $\rm Ti_{1.7}Sb$  phase and the  $\rm Ti_{1.2}Sb$  phase, were identified in the titanium-rich (> 50 at. % Ti) part of the system. The exact composition ranges of the phases have not been determined, and the formulae only indicate their approximate composition. Phase-analytical work using metallographic and X-ray powder methods and density measurements is still in progress and only a few observations will be mentioned here.

The  $Ti_{1.7}Sb$  phase. X-Ray powder photographs of different samples in the range around  $Ti_2Sb$  contained reflections from a new phase with composition somewhere in this concentration range. This phase was obtained by either of the melting methods outlined in the experimental part. By means of metallographic methods the approximate composition of this phase is found to be  $Ti_{1.7}Sb$ , which is in agreement with the composition estimated using the disappearing phase method on the powder photographs.

Single crystals of the Ti<sub>1.7</sub>Sb phase were obtained with bipyramidal shape in different melts. Oscillation and Weissenberg photographs showed that the crystals had orthorhombic symmetry. With the approximate unit cell dimen-

Table 6. Powder photograph	data of Ti <sub>1.7</sub> Sb taken	with $CrK$ -radiation, $a_i$	$_2$ and $\beta$ lines
	omitted.		-

$I_{ m obs}$	$\sin^2\!\Theta$	$egin{array}{c c} \sin^2\!\Theta  imes 10^4 & & & & & & & & & & & & & & & & & & &$		$\sin^2\!\Theta$	hkl			
1005	obs	calc	70700	2005	obs	cale	70100	
vw	954	953	211	w	4164	4166	142	
w	1032	1030	002	vw	4212	4209	323	
$\mathbf{m}$	1140	1137	121	vw	4309	4307	014	
$\mathbf{m}$	1262	1259	220	w	4396	4384	512	
vw	1348	1345	112	vw	4513	4516	233	
$\mathbf{m}$	1399	1398	301		5000	15293	441	
$\mathbf{m}$	1909	1909	122	w	5299	5311	620	
$\mathbf{v}\mathbf{w}$	1955	1950	031			1		
$\mathbf{w}$	2027	2026	400		F400	(5448	314	
$\mathbf{v}\mathbf{w}$	2076	2077	131	vw	5460	\\ 5453	143	
m	2152	2150	321		{	(5466	251	
$\operatorname{st}$	2287	(2284	401	w	5565	5569	621	
st	2281	2289	222	vw	5814	5811	034	
w	2338	2632	113	m	6237	6235	523	
$\mathbf{v}\mathbf{w}$	2476	2472	411	vw	6572	6563	105	
w	2508	2505	013	vw	6876	6276	603	
w	2781	2778	420	m	7180	7175	533	
vw	2842	2849	132	m	7326	7315	125	
$\mathbf{v}\mathbf{w}$	2920	2922	322	m	7479	7473	514	
***	3018	<b>(3009</b>	040	m	7988	7987	$\bf 722$	
m	3018	3012	213	w	8106	8105	800	
$\mathbf{w}$	3578	3576	223	vw	8272	8268	344	
vw	3724	3718	430	w	8551	8551	811	
$\mathbf{v}\mathbf{w}$	3808	3808	422	w	8572	8568	633	
vw	3862	3862	332	m	8796	8796	460	
vw	4025	${f 4009} {f 4093}$	$\begin{array}{c} 033 \\ 042 \end{array}$	w	8900	8897	552	

sions from the single crystal data the lattice constants of the Ti<sub>1.7</sub>Sb phase were determined from the powder photographs, cf. Table 6, to be:

$$a = 10.172 \text{ Å}, \qquad b = 8.346 \text{ Å}, \qquad c = 7.135 \text{ Å}.$$

The only systematic missing reflections were of the type hk0 absent when h = 2n + 1. The probable space groups are thus  $Pmma\ (D_{2h}^5)$ ,  $Pmc2_1\ (C_{2v}^2)$  and  $Pma2\ (C_{2v}^4)$ .

The Ti<sub>1.2</sub>Sb phase. Samples with composition in the range about Ti<sub>1.2</sub>Sb give very complicated powder photographs and the phases are not easily identified by X-ray methods. By metallographic technique, however, two phases with approximate composition Ti<sub>1.2</sub>Sb and TiSb could be identified. In the same way as the Ti<sub>1.7</sub>Sb phase, the Ti<sub>1.2</sub>Sb phase was obtained by either of the melting methods.

162

721

651

731

570

452

602

1670

1672

1775

1783

1790

(1846

1851

1678

1780

1854

vw

 $\mathbf{m}$ 

$I_{ m obs}$	$\sin^2\!\Theta$	× 104	$oxed{hkl} oxed{I_{\mathrm{obs}}}$		$\sin^2\Theta$	$\sin^2\Theta \times 10^4$	
1008	obs	calc	Tekt	1 obs	obs	calc	hkl
vw w	329 386	327 384	121 140	st	1039	$\begin{cases} 1031 \\ 1038 \end{cases}$	610 161
w	448	${448}\atop 452}$	400 330	w	1078	1043	$egin{array}{c} 222 \\ 132 \\ 261 \\ \end{array}$
w vw	489 536	484 537	311 420	vw	1128 1208	$\begin{cases} 1122 \\ 1201 \\ 1209 \end{cases}$	270 630
w w	589 611	584 608	150 340	vw	1310	$\begin{cases} 1308 \\ 1310 \end{cases}$	$\begin{array}{c} 630 \\ 621 \\ 242 \end{array}$
w	666	668 678	250 241			(1312)	$\begin{array}{c} 412 \\ 631 \end{array}$
w	725 826	$\begin{cases} 723 \\ 818 \\ 828 \end{cases}$	510 341 160	vw	1421	$\left\{egin{array}{l} 1422 \ 1426 \end{array} ight.$	$\begin{array}{c} 080 \\ 152 \end{array}$
w	882 895	878	251	vw vw	1514 1607	1510 1605	252 711

Table 7. Powder photograph data of low angle reflections of Ti<sub>1,2</sub>Sb taken with CuK-radiation,  $a_2$  and  $\beta$  lines omitted.

In most of the samples melted in the Nernst-Tammann furnace long needleshaped crystals were obtained. Oscillation and Weissenberg photographs showed the crystals to be orthorhombic with unit cell dimensions:

$$a = 14.55 \text{ Å}, \qquad b = 16.34 \text{ Å}, \qquad c = 5.31 \text{ Å}.$$

Missing reflections were: 0kl when k = 2n + 1 and h0l when h = 2n + 1. This is characteristic of the space groups  $Pbam~(D_{2h}^9)~$  and  $Pba2~(C_{2v}^8)$ . In Table 7 are listed the lines found in the low angle region on a powder photograph of the Ti<sub>1,2</sub>Sb sample.

The TiSb phase. X-Ray powder photographs of the TiSb phase were indexed on the basis of a hexagonal unit cell. The redetermination of the lattice constants gave the result:

$$a = 4.115 \text{ Å}, \qquad c = 6.264 \text{ Å}, \qquad c/a = 1.560$$

Whether the difference between the values given here and the results by Nowotny and Pesl<sup>1</sup> ( $a = 4.07_0$  Å;  $c = 6.30_6$  Å) reflects the difference in thermal history of the alloys, or indicates an extended homogeneity range, or is due to impurities is not clear at the moment.

895

1006

1017

m

m

892

1000

1004

1009

1010

1014

1018

112

521

450

600

061

441

#### DISCUSSION

In the  $\beta$ -W structure of the cubic  $Ti_3Sb$  phase, cf. Fig. 2, each titanium atom is coordinated to two titanium atoms in linear chains at a distance a/2, to eight more titanium atoms at a 22.5 % larger distance, and to four antimony atoms in a tetrahedral arrangement. Each antimony atom is coordinated to twelve titanium atoms at the corners of an icosahedron. The interatomic distances between nearest neighbours in cubic  $Ti_3Sb$  are listed in Table 8.

Table 8. Interatomic distances in cubic Ti<sub>3</sub>Sb (Å).

Ti — 2 Ti : 2.609 — 8 Ti : 3.195 — 4 Sb : 2.917 Sb — 12 Ti : 2.917

The shortest Ti—Ti distance is definitely shorter than in metallic titanium. The shortest interatomic distances are 2.896 Å for  $\alpha$ -Ti according to Clark 9 and 2.843 Å for  $\beta$ -Ti according to Levinger <sup>10</sup>. In many of the titanium compounds with  $\beta$ -W structure, listed by Geller <sup>11</sup>, even shorter Ti—Ti distances than the one found here are observed.

The observed interatomic distances can also be compared with the values predicted from the metallic radii 12. If antimony retains one of its usual valences, 1.56 or 3, the twelve Ti-Sb bonds should have bond numbers 0.13 or 0.25. Using the single bond radius value of 1.324 Å for Ti and 1.391 Å for Sb, with an electronegativity correction of 0.008 Å, the expected Ti-Sb bond distances for the two different valence states are 3.239 Å and 3.068 Å, respectively. The observed value 2.917 Å agrees better with that calculated for a valence 3 for Sb than for the usual metallic valence 1.56, but even this value is considerably larger than the observed one. To account for the discrepancy one can either accept a set of special  $\beta$ -W radii (cf. Geller 11,13 and Pauling 14), or assume a higher bond number for the Ti-Sb bonds 14. Higher bond numbers would obtain, according to Pauling 12, if one electron were transferred from the hyperelectronic Sb atom to the hypoelectronic Ti atom, thus increasing the valences by one. The calculated Ti—Sb bond distance for this valence, 2.993 Å is still appreciably larger than the observed value. Using the effective metallic radius for Ti given by Pauling 14 for the β-W structure, a shorter Ti-Sb bond distance value of 2.940 Å for the same bond number 0.33 is calculated. For the ten Ti-Ti bonds with a bond number of 0.20, and using the  $\beta$ -W radius value for singlebonded Ti of 1.271 Å, the calculated bond distance is 3.058 Å, or rather close to the weighed mean (3.078 Å) of the observed values.

In order to derive the structure of tetragonal  $\text{Ti}_3\text{Sb}$  from the  $\beta$ -W type structure of cubic  $\text{Ti}_3\text{Sb}$ , displacements of some of the atoms must be performed. The main difference between the two structures consists in a movement of the Ti I atoms as indicated by the arrows in Fig. 6 and a translation of these atoms of magnitude c/4 along the c-axis. Moreover, a characteristic difference is produced by the change from cubic to tetragonal symmetry, cf. Figs. 5 and 6.

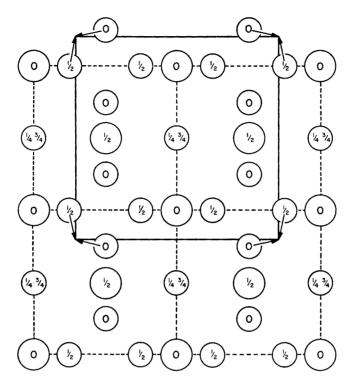


Fig. 6. Projection of four unit cells of cubic Ti<sub>3</sub>Sb in relation to the 001 projection of the primitive unit cell of tetragonal Ti<sub>3</sub>Sb.

Chains of Ti atoms, characteristic of the cubic Ti<sub>3</sub>Sb structure are also present in the tetragonal structure. In the tetragonal structure Ti I and Ti II form chains only along the c-axis, while perpendicular to it the chains are broken due to the relative displacement of the Ti I atoms and by the change to tetragonal symmetry. The configuration of the eight Ti III atoms around each Ti I is found to be that of a slightly distorted square antiprism. The Ti II atoms are coordinated to eight Ti III atoms at the corners of an even more distorted rectangular antiprism and to four antimony atoms at the corners of a tetrahedron. Each Ti III atom is surrounded by eleven titanium atoms and four antimony atoms, cf. Fig. 5. Each antimony atom is coordinated to ten titanium atoms and two antimony atoms at the corners of an irregular icosahedron.

The interatomic distances between nearest neighbours in tetragonal Ti<sub>3</sub>Sb calculated for the final parameters are listed in Table 9. The errors indicated are standard deviations calculated according to Cruickshank <sup>15</sup>. No comparison between observed and calculated bond distances is carried out for tetragonal Ti<sub>3</sub>Sb, since even in the case of cubic Ti<sub>3</sub>Sb the basis for the calculation is uncertain.

Table 9. Interatomic distances in tetragonal Ti<sub>3</sub>Sb (Å).

```
2 Ti
                       I : 2.632
         — 12 Ti
                     III : 2.794 \pm 0.011
             2 Ti
8 Ti
                      II : 2.632
                            3.285 \pm 0.011
                     III:
             4 Sb
                         : 2.755 \pm 0.003
Ti III
             2 Ti
                       I : 2.794 \pm 0.011
              2 Ti
                      II : 3.285 \pm 0.011
              1 Ti
2 Ti
                     III : 2.970 \pm 0.021
                     III : 3.082 \pm 0.021
             2 Ti
2 Ti
                     III : 3.407 \pm 0.021
                     III : 3.487 \pm 0.021
              1 Sb
                         : 2.780 \pm 0.012
              1 Sb
                            2.790 \pm 0.012
                         :
              2 Sb
                            3.029 \pm 0.012
              2 Ti
2 Ti
                      \mathbf{II}:
Sb
                            2.755 \pm 0.003
                     III : 2.780 \pm 0.012
              2 Ti
                     III : 2.790 \pm 0.012
              4 Ti
                     III : 3.029 \pm 0.012
                          : 3.671 + 0.004
```

In order to study relative stabilities of the two crystalline phases, samples first melted in the arc furnace were remelted in the Nernst-Tammann furnace, and *vice versa*. By this treatment, the cubic crystals were changed into tetragonal crystals, while the tetragonal crystals remained unchanged after melting in the arc furnace. Thus, the transformation is apparently irreversible.

It is known in other instances that small amounts of impurities can stabilize more or less foreign structures, e.g. a pseudobrookite structure for Ti<sub>3</sub>O<sub>5</sub> is found already at an iron concentration of 0.23 weight % Fe (x = 0.03 in (Ti<sub>1-x</sub>Fe<sub>x</sub>)<sub>3</sub>O<sub>5</sub>) according to Andersson et al. 16 Spectrographic analysis of samples containing the tetragonal Ti<sub>3</sub>Sb phase showed the presence of about 3/10 % Al and 3/10 % Si as impurities. Samples of the cubic Ti<sub>2</sub>Sb phase were, however, only found to contain the same amounts of impurities as the titanium and antimony used. On melting samples with composition Ti<sub>3</sub>Sb in the arc furnace together with different small amounts of Al, Si, Al<sub>2</sub>O<sub>3</sub> and SiO<sub>2</sub> they all still gave X-ray photographs characteristic of the cubic Ti<sub>3</sub>Sb phase. The presence of impurities like oxygen and nitrogen might perhaps explain the stability of the tetragonal structure. Thus, cubic Ti<sub>3</sub>Sb is obtained by remelting tetragonal Ti<sub>3</sub>Sb in an arc furnace with small amounts of pure titanium. The explanation of this as due to differences in composition of the two phases, seems unreasonable since, for example, the alloy Ti<sub>3.5</sub>Sb remains tetragonal after remelting in the arc furnace. Indeed, the pure titanium might act as an absorber of the oxygen and nitrogen present in the alloy.

Including the four phases reported here eight phases have been described in the titanium-antimony system. Their characteristic crystallographic properties are listed in Table 10. In view of the present discussion about the relative stability of cubic and tetragonal Ti<sub>3</sub>Sb, and as the present study could not confirm the existence of the Ti<sub>4</sub>Sb phase and the Ti<sub>2.5</sub>Sb phase reported by Nowotny and coworkers <sup>1-3</sup>, the possibility that these phases also might be

Phase	d (g cm <sup>-3</sup> )	a (Å)	b (Å)	c (Å)	c/a	Structure type	Reference
Ti <sub>4</sub> Sb	_	5.958		4.807	0.807	HexNi <sub>3</sub> Sn	Nowotny et al.2
Ti₃Sb	6.137	5.2186	_		_	CubβW	Present
Ti <sub>3</sub> Sb	-	5.217			_	CubβW	Matthias et al.6
Ti₃Sb	6.070	10.465		5.2639	0.508	Tetr.	Present
$Ti_{2.5}Sb$	5.90	4.01		14.5	3.63	TetrCu <sub>2</sub> Sb	Auer-Welsbach
						_	et al.3
$Ti_{1.7}Sb$	_	10.172	8.348	7.135	_	Orthor.	Present
$Ti_{1,2}Sb$		14.55	16.34	5.31	_	Orhtor.	Present
TiSb	-	$4.07_{0}$	_	6.30	1.55	HexNiAs	Nowotny and Pesl 1
TiSb	_	4.115		6.264	1.560	Hex.	Present
TiSb <sub>2</sub>		$6.66_{6}$		5.81,	$0.87_{3}$	TetrCuAl <sub>2</sub>	Nowotny and Pesl <sup>1</sup>

Table 10. Pycnometric densities and crystallographic data for phases reported in the titanium-antimony system.

stabilized by impurities should not be overlooked. With attention focused on the Ti<sub>4</sub>Sb phase it should be noted that a lattice constant variation of only 2 to 4 % between the isostructural phases Ti<sub>4</sub>Sb, Ti<sub>4</sub>Pb (Nowotny and Pesl 17) and Ti<sub>3</sub>Ål \* (Anderko et al. 18) is much smaller than the 10 to 23 % difference expected due to the difference in metallic radii of Sb, Pb and Al. This point is perhaps worth a further study in view of the observed reaction between Ti and Al<sub>2</sub>O<sub>3</sub>.

Acknowledgments. The authors are greatly indebted to Professor Haakon Haraldsen for his kind interest in this study and for placing laboratory facilities at their disposal. They also wish to express their thanks to siviling. Hallstein Kjøllesdal, Sentralinstitutt for Industriell Forskning, for carrying out the meltings in the arc furnace.

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<sup>\*</sup> The Ti<sub>2</sub>Al phase reported by Ence and Margolin <sup>19</sup> is probably belonging to the homogeneity range of the Ti, Al phase.

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Received January 12, 1962.