

Fig. 2. Dosage curves for vanillin-blue:

a. Diluted with 70 % sulphuric acid 0 ——— 0
b. Diluted with 96 % sthanol + ——— +

tion with ethanol, and  $(28.4 \pm 0.9) \times 10^8$  units/mole, for sulphuric acid dilution. The advantage in diluting with ethanol is the maximal absorption at  $600 \text{ m}\mu$ . The interference with other coloured substances formed, i. e., the red substance in the supernatant liquid will be less in ethanol. On the other hand the alcoholic solutions are less stable. A point to be noted in the practical application of the method is that both an aged vanillin-sulphuric acid solution and a freshly prepared one will give the same dosage curve.

For the use of vanillin-blue as a peroxide reagent, it is unnecessary to know its formula or the reaction mechanism. The amount of vanillin exceeds greatly that of the peroxides (ranging from 1:10 to 1:104). It is doubtful whether the vanillin-blue formed and precipitated has a constant composition. However, an analysis was made, giving C 68.2, H 5.46, and O 24.3 (1 % ash).

Campbell and Coppinger 2 added t-butyl hydroperoxide to 2,6-di-t-butyl-p-cresol, and obtained a substance described as 1-methyl - 1 - t - butylperoxy - 3,5 - di - t-butylcyclohexadienone-4. With this type of peroxide (which have an absorption maximum at 234 m $\mu$ ) vanillin-blue seems to have little in common. We exclude a compound with the t-BuOO-group directly attached to the aromatic ring as such compounds will not be stable. In view of the small amounts of peroxide needed for the formation of colour and the unknown be-

haviour of less stable peroxides such as the methylhydroperoxide, no formula for vanil-lin-blue is advanced.

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Received March 19, 1955.

## 4s-Electrons in the First Transition Group Complexes CHR. KLIXBULL JØRGENSEN

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When the crystal field theory is applied  $^{1,2}$  to the absorption spectra of hexaaquo ions and other simple complexes of the first transition group, small discrepancies ( $\sim 5$  %) occur between the calculated and observed wavenumbers of the band maxima. Mainly based on magnetic evidence, Owen  $^3$  has interpreted these effects by  $\sim 20$  % intermixing of  $\sigma$ -electrons from the ligands with the  $\gamma_3$ -electrons (the highest energy state of d-electrons in an octahedral complex while the other is  $\gamma_5$ ). In the present note attention will be drawn to the electron configuration  $3d^{n-1}$  4s, which in the free, divalent ion is situated  $^4$   $\sim 50~000~{\rm cm}^{-1}$  over the groundstate, due to  $3d^n$ .

Since these two electron configurations have the same parity (the sum of the 1values is even) they can directly intermix. The author 5 even maintains that the crystal field induces an intermixing with 3dn-14p and electron transfer states, explaining the observed transitions which would else be forbidden by Laporte's rule. The former interaction will inter alia depress the energy of the levels with the same  $\Gamma_n$  as the lowest level of the excited  $3d^{n-1}$  4s configuration. Thus, in  $d^3$ -systems, the levels  ${}^4\Gamma_4$  will be depressed by the ground-state  $\gamma_5{}^2\gamma_1$  of  $3d^24s$ . Now, Mr. C. E. Schäffer has kindly pointed out to me that the observed energy difference,  ${}^4\Gamma_4 - {}^4\Gamma_5$ , between the two strong bands of chromium (III) complexes decreases by increasing crystal field strength, while it should 6 be slightly increasing or constant, if no interaction appeared. Also V<sup>++</sup> shows this effect if  ${}^4\Gamma_5$  is placed at 12 200 cm<sup>-1</sup>, contrary to Owen 3.

In di-systems as Mn++, Orgel discussed a systematical depression of <sup>4</sup>G and other excited terms in the hexaaquo ion, compared with the free ion 4. This might also be explained by interactions especially with the low '\(\Gamma\_a\) and '\(\Gamma\_4\) of 3d'4s. With the Cary spectrophotometer a very broad band  $(\delta \sim 1.500 \text{ cm}^{-1})$  is observed of 2 M MnSO. in H<sub>2</sub>O at 32 600 cm<sup>-1</sup>. This is undoubtedly due to  ${}^4\Gamma_4$ , and the band at 29 700 cm<sup>-1</sup> to  ${}^4\Gamma_a$ , as pointed out by Orgel<sup>2</sup>. At 38 500 and at 40 600 cm<sup>-1</sup>, two bands are observed, probably due to  ${}^4\Gamma_4$  and  ${}^4\Gamma_5$  of the term \*F, which should be displaced upwards by perturbations <sup>3</sup>. <sup>4</sup>T<sub>2</sub> can be assigned to one of the two shoulders observed at 35 900 and 37 000 cm<sup>-1</sup>. The three latter levels show definitely the decrease in energy of \*F. due to interactions with states of other electron configurations, e. g. 3d4s. All the quartet levels of 3d5 in Mn++ seem now detected.

In d<sup>8</sup>-systems as Ni<sup>++</sup>, the lowest levels of  $3d^74s(\gamma_5^6\gamma_3\gamma_1)$  are  $^3\Gamma_3$  and  $^1\Gamma_3$ . The latter level will depress the lowest singlet level 7 of nickel(II) complexes. Without this interaction the transition  ${}^3\Gamma_2 - {}^1\Gamma_3$  would correspond to the weak band  ${}^2$ , at 18 350 cm<sup>-1</sup> of Ni(H<sub>2</sub>O)<sub>6</sub>++. But the "extra" band at 15 400 cm<sup>-1</sup> has been assigned 5 to this transition. In a paper to be published the energy of 1/2 in a wide range of nickel(II) complexes will be studied. Table 1 shows the observed maxima in cm<sup>-1</sup> of complexes, which are not all of cubic symmetry (e. g. the ethylenediaminetetraacetate or the solutions in concentrated acids). The parentheses indicate cases of very strong intermixing, where the distinction between 1/3 and the triplet state has no physical significance 8. This results in nearly equal intensities and halfwidths, while  ${}^{1}\Gamma_{3}$  ordinarily is weak and narrow.

While the band 'should be quite constant ~ 18 000 cm<sup>-1</sup> as an intermixing of 'D and 'G found by Shenstone', the observed values decrease considerably with increased crystal field strength. In tetragonal complexes the interaction will be even more pronounced. Thus, the energy difference, necessary to make a planar complex diamagnetic, will be rather 11 000 than 18 000 cm<sup>-1</sup> (cf. eq. 29 of Ref. 7). This agrees better with the observed tendency of changing groundstate.

The ethanol solvate of nickel(II)bis-(acetylacetonate), which probably is octahedral <sup>10</sup>, has a shoulder at 13 000 cm<sup>-1</sup> besides the stronger triplet bands at 9 100 and 15 550 cm<sup>-1</sup>. If nickel nitrate in 60 %

Table 1.

	$^{\scriptscriptstyle 1} arGamma_{f 3}(D)$	${}^{3} {\varGamma}_{5}(F)$	$^3 \Gamma_4(F)$	${}^3\Gamma_4(P)$
free	14 000	0	0	16 900
H <sub>2</sub> SO <sub>4</sub>	14 800		$12\ 200$	23 350
H <sub>3</sub> PO <sub>4</sub>	14 900		13 150	24 500
$(H_2O)_6$	$(13\ 500)$	8 500	(15 400)	25 300
enta	12 700	10 100	17 000	26 200
glycine <sub>a</sub>	13 100	10 100	16 600	27 600
(NH <sub>a</sub> )	13 150	10 750	17 500	28 200
en <sub>s</sub>	12 400	11 200	18 350	29 000
aa'-dipa	$(11\ 500)$	(12650)	19 200	
o-phen <sub>3</sub>	(11 550)	(12 700)	19 300	_

glycerol is cooled by liquid air, a sharp band in the far red can be distinguished in a spectroscope, corresponding to a later place in the Table. The non-diagonal elements  $^8$  in the matrices between  $^1D$  and  $^8F$ , due to (L, S) coupling, are half the smallest distances between the two bands, when the levels are crossing,  $\sim 800 \text{ cm}^{-1}$  in the table. If the non-diagonal elements between  $3d^n$  and  $3d^{n-1}4s$  are of the order of magnitude of  $10~000~\text{cm}^{-1}$ , then the intermixing in squares of the wavefunctions will be about 4~%, and the energy decreases of the lowest levels  $\sim 2~000~\text{cm}^{-1}$ , agreeing well with the observed effects.

Transitions between the configurations 3dn and 3dn-14s can next be sought for in the absorption spectra. Orgel 2 pointed out that copper(I) complexes do not show such bands in the wavenumber range  $21\ 900-26\ 300\ {\rm cm^{-1}}$ , where levels of  $^3D$  and  $^1D$  of  $3d^9$  4s are distributed in the free ion 4. Solutions of Cu(NH<sub>3</sub>)<sub>2</sub>+ in 0.2 M NH<sub>4</sub>ClO<sub>4</sub> and 1 M NH<sub>2</sub> show on the Cary absorption above 35 000 cm<sup>-1</sup>. If this is identified as these transitions, giving not much higher intensities than the usual crystal field spectra, it is seen that the 4s-electron has considerably higher energy in the crystal field than in the free ion. This is formally connected with the different values 11 of the crystal field parameter  $B_0$  in different configurations. Orgel 12 has reviewed the "electron transspectra and agrees with Dainton that some bands with  $\varepsilon \sim 100$  of divalent ions in the ultraviolet are due to transfer of electrons from the central ion to the ligands, the opposite way of the ordinary bands (with  $\varepsilon \sim 5\,000$ ) found in oxidizing metal ions such as Fe+3, Cu+2, Ir+4, Pu+4, etc. The former type of band, observed in V++ at 33 000 cm<sup>-1</sup> and in Cr++ at 40 000  ${\rm cm}^{-1}$  may be ascribed to the  $3d^{n-1}4s$  states. Since  ${}^4F$  and  ${}^5F$  of these configurations are

situated at 44 000 and 49 000 cm<sup>-1</sup> in the free ions 4, the crystal field splittings of the excited terms must here be considerable in order to explain the low wavenumbers observed. It might be argued that 'S of 3d<sup>5</sup>4s in Fe++ at 30 000 cm<sup>-1</sup> should give even lower wavenumbers. But for more than five d-electrons, the highest multiplicity of 3d<sup>n-1</sup>4s gives spin-forbidden bands, and first 'S at 41 000 and 'G at 63 000 cm<sup>-1</sup> will give ordinary intensities in Fe++, where a band '2 is observed ~40 000 cm<sup>-1</sup>.

Orgel and Owen investigate the possibility of covalent bonding, i. e. molecular orbitals being occupied by electrons from both central ion and ligands. This is undoubtedly the case of ligands with considerable electron affinity, as CO, CN-, NO+, aromatic amines, PCl<sub>3</sub>, trialkylphosphines, but it is not easily decided in the case of ordinary ligands (H<sub>2</sub>O, NH<sub>3</sub>, etc.) where the most conspicuous effect <sup>2,3</sup> on the absorption spectra is only an increased energy difference  $(E_1-E_2)$  between  $\gamma_3$  and  $\gamma_5$ electrons, analogous to the crystal field influence. It is interesting that  $(E_1-E_2)$ is nearly constant ~20 000 cm<sup>-1</sup> in trivalent hexaaquo ions, while it is  $\sim 10~000$  cm<sup>-1</sup> in divalent ions, and here decreasing <sup>1</sup> with the atomic number as implied from the theory 13. This can only be explained by considerably smaller distances to the effective negative charges of the ligands of the trivalent ions than in the divalent ions, if the covalent hypothesis is not accepted. It must be remarked that the parameters in the crystal field model of Ilse and Hart-mann 13 have no quantitative physical significance. E. g., the hydrogen-like 3d wavefunction with the effective charge Z = 4 has its maximum at a distance 1.18 A from the nucleus, while the radius of Ti+3 is assumed to be 0.8 A. Some problems related to effective quantum numbers will be discussed elsewhere 14. The transitions between different configurations,  $5f^n \rightarrow 5f^{n-1}6d$  in the actinide ions, are also known from absorption spectra 15.

A valuable implication from the theory of molecular orbitals  $^{2,3}$  is that the strong electron transfer bands are due to transitions from the odd  $\gamma_4$ -states. Hartmann  $^{16}$  pointed out that the series of energy of the molecular orbitals in an octahedral complex should be:

 $\gamma_1$ ,  $\gamma_4$ ,  $\gamma_3$ ,  $\gamma_5(d)$ ,  $\gamma_3(d)$ ,  $\gamma_1(s)$ ,  $\gamma_4(p)$ ...

If the lowest  $\gamma_3$  had a much higher energy than the lowest  $\gamma_4$ , weak bands would be

found in the electron transfer spectra at lower wavenumbers than the strong bands.

Acknowledgments. I am much indebted to Professor Jannik Bjerrum for interesting discussions. Further, I thank Dr. L. E. Orgel for the opportunity to see the manuscripts of several new papers.

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Received March 24, 1955.

## Dihydro-thionaphthene-2- and -3-carboxylic Acids

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The acids have been prepared in connection with current work on optically active plant growth substances. They are easily obtained by hydrogenation of the corresponding thionaphthene-carboxylic acids using sodium amalgam.

Acta Chem. Scand. 9 (1955) No. 4